

MOMENT MAPS AND THE GEOMETRY OF THREE - FORMS

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Plan

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Multi-moment maps
Topological existence

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Lie algebras
Theory of orbits

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I, GEOMETRY

Three-forms on vector spaces

$$V = \mathbb{R}^n$$

$$c \in \wedge^3 V^*$$

Definition: c is fully non-degenerate if

$$c(v, w, \cdot) \neq 0$$

$\forall$ linearly indep. $v, w \in V$

Theorem (Massey, 1983): c fully non-degenerate $\Rightarrow n = 3$ or 7 .

Definition: c is (weakly) non-degenerate if

$$c(v, \cdot, \cdot) \neq 0 \quad \forall v \in V, v \neq 0.$$

Proposition: c non-degenerate $\Rightarrow n \neq 1, 2, 4$

Definition (Hitchin): c is stable if $GL(V) \cdot c \subseteq \wedge^3 V^*$ is open.

Proposition: c stable $\Rightarrow n = 3, 4, 5, 6, 7, 8$

M^n a manifold of dimension n .

Definition: a strong geometry is a pair (M, c) , $c \in \Omega^3(M)$ such that

$$dc = 0$$

It is 2-plectic (Baez, Hoffnung, Rogers)
if $c_m \in \Lambda^3 T_m^* M$ is non-degenerate $\forall m \in M$

Examples

(1) Extended phase space

$$M = \Lambda^2 T^* N$$

$$c = db$$

b the tautological 2-form

(2) $M = S^1 \times N$ (N, ω) symplectic
 $c = dt \wedge \omega$

(3) $M = G$ a compact Lie group
 $c(x, y, z) = \langle [\theta^L(x), \theta^L(y)], \theta^L(z) \rangle$

$\langle, \rangle$ a biinvariant inner product
on $\mathfrak{g} = T_e G$

$$\theta^L(x_g) = (L_{g^{-1}})_* X \in \mathfrak{g}$$

$L_g = gh$, left-multiplication.

2-plectic $\iff G$ semi-simple

(4) strong KT geometry

(M, I, g) a Hermitian manifold

g - Riemannian metric

I - integrable complex structure

$$g(IX, IY) = g(X, Y) \quad \forall X, Y$$

$$\omega_I(X, Y) := g(IX, Y), \quad \omega_I \in \Omega^2(M)$$

This is Kähler if $d\omega_I = 0$

It is strong KT if $\partial_I \bar{\partial}_I \omega_I = 0$
i.e. $dI d\omega_I = 0$

$$c = -I d\omega_I \in \Omega^3(M)$$

is then closed, but can be degenerate.

$$\nabla^B = \nabla^K + \frac{1}{2}c$$

is the Bismut connection

(5) strong HKT geometry

(M, g, I, J, K)

$$IJ = K = -JI$$

$$c = -I d\omega_I = -J d\omega_J = -K d\omega_K$$

$\Rightarrow I, J, K$ integrable

$+dc = 0$ is strong HKT

E.g. (i) hyper-Kähler manifolds

(ii) most compact G with $\dim G = 4r$
(Joyce)

Open problem: find other compact simply-connected examples.

Multi-moment maps

Suppose (M, c) is a ~~class~~ strong geometry and G is a connected group of symmetries i.e.

$$L_X c = 0 \quad \forall \underline{X} \in \mathfrak{g}$$

Cartan's formula

$$\begin{aligned} 0 &= L_X c = d(X \lrcorner c) + X \lrcorner \cancel{dc}^0 \\ &= d(X \lrcorner c) \quad \text{for } \underline{X} \in \mathfrak{g}. \end{aligned}$$

For $\underline{Y} \in \mathfrak{g}$,

$$\begin{aligned} L_Y (X \lrcorner c) &= (L_Y X) \lrcorner c = -[X, Y] \lrcorner c \\ &= d(Y \lrcorner (X \lrcorner c)) + Y \lrcorner \cancel{d(X \lrcorner c)}^0 = d((X \wedge Y) \lrcorner c) \end{aligned}$$

Thus for $\underline{p} = \sum \underline{X}_i \wedge \underline{Y}_i \in \wedge^2 \mathfrak{g}$

$\underline{p} \lrcorner c \in \Omega^1(M)$ is closed

$$\iff L(\underline{p}) = 0 \quad \text{where } L = [\cdot, \cdot]$$

Definition: the Lie kernel of $\mathfrak{g}$ is

$$\mathcal{P}_{\mathfrak{g}} = \ker(L = [\cdot, \cdot] : \wedge^2 \mathfrak{g} \longrightarrow \mathfrak{g})$$

Examples: $\mathfrak{g}$ Abelian $\Rightarrow \mathcal{P}_{\mathfrak{g}} = \wedge^2 \mathfrak{g}$

$\mathfrak{g}$ semi-simple $\Rightarrow \wedge^2 \mathfrak{g} = \mathfrak{g} \oplus \mathcal{P}_{\mathfrak{g}}$

and (hard) $\mathcal{P}_{\mathfrak{g}}$ is an irreducible module when $\mathfrak{g}$ is simple

Definition: $\nu: M \longrightarrow \mathcal{P}_g^*$

is a multi-moment map
if it is equivariant and

$$d\langle \nu, p \rangle = p \lrcorner c$$

$$\forall p \in \mathcal{P}_g$$

Example: $M = \Lambda^2 T^*N$ $c = db$

$$G \subseteq \text{Diff}(N)$$

$$\langle \nu, p \rangle = b(p)$$

Theorem: Suppose $b_1(M) = 0$.

Then ν exists if

either (a) G is compact

or (b) M is compact with

a G -invariant volume form

Example: $M = S^1 \times N$ (N, ω) symplectic

If H acts in a Hamiltonian way on N , moment

$$\text{map } \mu: N \longrightarrow \mathfrak{h}^*$$

then let

$$G = S^1 \times H$$

$$\mathcal{P}_G \cong \mathbb{R} \wedge \mathfrak{h} + \mathcal{P}_{\mathfrak{h}}$$

$$\left\{ \begin{array}{ll} \langle \nu, p \rangle := 0 & \text{for } p \in \mathcal{P}_{\mathfrak{h}} \\ \langle \nu, T_x X \rangle := \langle \mu, X \rangle & \text{for } x \in N \end{array} \right.$$

gives a multi-moment map.

MULTI-MOMENT MAPS

II, Algebra

Recall: • (M, c) strong geometry $c \in \Omega^3(M)$
 $dc = 0$

• G a connected group of symmetries
Lie kernel $\mathcal{P}_g = \ker (L = [\cdot, \cdot]: \Lambda^2 \mathfrak{g} \rightarrow \mathfrak{g})$

• a multi-moment map is

$$\nu: M \longrightarrow \mathcal{P}_g^* \quad \text{equivariant}$$

with

$$d\langle \nu, \varphi \rangle = \varphi \lrcorner c \quad \forall \varphi \in \mathcal{P}_g^*$$

Today: • algebraic existence & uniqueness
• structure theory
• orbit theory

Have

$$0 \rightarrow \mathcal{P}_g \rightarrow \Lambda^2 \mathfrak{g} \xrightarrow{L} \mathfrak{g} \rightarrow \mathfrak{g}/\mathfrak{g}' \rightarrow 0$$

$\mathfrak{g}' = [\mathfrak{g}, \mathfrak{g}]$ the derived algebra

so $\dim \mathcal{P}_g = \dim \Lambda^2 \mathfrak{g} - \dim \mathfrak{g}'$

Dually

$$0 \rightarrow \text{Ann}(\mathfrak{g}') \rightarrow \mathfrak{g}^* \xrightarrow{d} \Lambda^2 \mathfrak{g}^* \xrightarrow{\pi} \mathcal{P}_g^* \rightarrow 0$$

so $\mathcal{P}_g^* \cong \Lambda^2 \mathfrak{g}^* / B^1(\mathfrak{g})$

Theorem:

(a) if $b_2(g) = 0$, then v is unique
if it exists

(b) if $b_2(g) = 0 = b_3(g)$, then v exists
and is unique

Proof: $\mathcal{P}_g^* \cong \Lambda^2 g^* / B'(g)$

$$\text{ind} = B'(g) \subset Z^2(g) = \ker d$$

$$\Rightarrow d_p : \mathcal{P}_g^* \longrightarrow \Lambda^2 g^*$$

$$d_p(\beta + B'(g)) = d\beta \quad \text{well-defined,}$$

G -equivariant

$$\ker d_p = \ker d / B'(g) \quad \text{im } d_p = \text{im } d \subset Z^2(g)$$

$$(i) \quad d_p \text{ injective} \iff b_2(g) = 0$$

$$(ii) \quad \text{im } d_p = Z^2(g) \iff b_3(g) = 0$$

Define

$$\begin{aligned} \Psi : \mathcal{M} &\longrightarrow Z^2(g) \\ \langle \Psi(m), \underline{x} \wedge \underline{y} \wedge \underline{z} \rangle &= \underline{c}_m(x, y, z) \end{aligned}$$

When $b_2(g) = 0$,

v is a multi-moment map

$$\iff d_p v = \Psi$$

□

Definition: $\mathfrak{g}$ (or $\mathfrak{g}$) is
 (cohomologically) $(2, 3)$ -trivial
 if $b_2(\mathfrak{g}) = 0 = b_3(\mathfrak{g})$

Rough classification of Lie algebras:

(1) $\mathfrak{g}$ solvable if $\mathfrak{g}^k = 0$ for some $k \geq 1$,
 $\mathfrak{g}' = \mathfrak{g}' = [\mathfrak{g}, \mathfrak{g}]$, $\mathfrak{g}^{k+1} = [\mathfrak{g}^k, \mathfrak{g}^k]$

(2) $\mathfrak{g}$ nilpotent, if $\mathfrak{g}^{(k)} = 0$ for some $k \geq 1$,
 $\mathfrak{g}^{(1)} = \mathfrak{g}'$, $\mathfrak{g}^{(k+1)} = [\mathfrak{g}, \mathfrak{g}^{(k)}]$
 "upper triangular matrices"
 "strictly upper triangular"

(3) $\mathfrak{g}$ solvable $\Rightarrow \mathfrak{g}'$ nilpotent

(4) $\mathfrak{g}$ simple if it has no non-trivial
 proper ideals

(5) $\mathfrak{g}$ semi-simple if it is a direct
 sum of simple ideals

(6) $\text{rad}(\mathfrak{g})$ is the maximal solvable
 ideal of $\mathfrak{g}$

(7) $0 \rightarrow \text{rad}(\mathfrak{g}) \rightarrow \mathfrak{g} \rightarrow \mathfrak{g}_{ss} \rightarrow 0$
 $\mathfrak{g}_{ss}$ semi-simple

(8) $\mathfrak{g}$ semi-simple $\Rightarrow b_3(\mathfrak{g}) = \# \text{ simple summands}$

Theorem: $\mathfrak{g}$ is $(2,3)$ -trivial

$\iff \left\{ \begin{array}{l} \mathfrak{g} \text{ is solvable} \\ \text{(not a product)} \end{array} \right.$

$$\left\{ \begin{array}{l} b_1(\mathfrak{g}) = \dim(\mathfrak{g}/\mathfrak{g}') = 1 \\ \& \quad H^i(\mathfrak{g}')^{\mathfrak{g}} = 0 \quad i=1,2,3 \end{array} \right.$$

Example: $\mathfrak{g}$ $(2,3)$ -trivial, $\dim \mathfrak{g} = 3$

$$\Rightarrow \mathfrak{g}' = \mathbb{R}^2 \quad \mathfrak{g} = \mathbb{R}X + \mathfrak{g}' \quad H^1(\mathfrak{g}') = \mathbb{R}^2, \quad H^2(\mathfrak{g}') = \mathbb{R}$$

$$\quad \quad \quad = \wedge^1 \mathfrak{g}'^* \quad \quad \quad = \wedge^2 \mathfrak{g}'^*$$

$$\begin{array}{ll} H^1(\mathfrak{g}')^X = 0 & \Rightarrow X|_{\mathfrak{g}'} \text{ acts invertibly} \\ H^2(\mathfrak{g}')^X = 0 & \Rightarrow \text{Tr } X|_{\mathfrak{g}'} \neq 0 \end{array}$$

$$\Rightarrow X|_{\mathfrak{g}'} = \begin{cases} \begin{pmatrix} \lambda & 1 \\ 0 & \lambda \end{pmatrix} & \lambda \neq 0 \\ \begin{pmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{pmatrix} & \begin{array}{l} \lambda_1 \neq 0 \\ \lambda_2 \neq \lambda_1 \end{array} \\ \begin{pmatrix} \lambda & \mu \\ -\mu & \lambda \end{pmatrix} & \begin{array}{l} \lambda \neq 0 \\ \mu \neq 0 \end{array} \end{cases}$$

$$\text{Scale } X \rightarrow \begin{cases} \lambda = 1 \\ \lambda_1 = 1, -1 < \lambda_2 \leq 1 \\ \mu = 1, \lambda > 0 \end{cases}$$

$$\mathfrak{g} = \begin{cases} \mathfrak{f}_3 = (0, 21+31, 31) \\ \mathfrak{f}_{3,\lambda_2} = (0, 21, \lambda_2 \cdot 31) \\ \mathfrak{f}'_{3,\lambda} = (0, \lambda \cdot 21+31, -21+\lambda \cdot 31) \end{cases}$$

Example: if $\mathfrak{g}'$ has a positive grading

$$\mathfrak{g}' = \mathfrak{k}_1 + \dots + \mathfrak{k}_r$$

$$[\mathfrak{k}_i, \mathfrak{k}_j] \subseteq \mathfrak{k}_{i+j} \quad \forall i, j$$

Let $X|_{\mathfrak{k}_i}$ be multiplication by i

Then $(\wedge^k \mathfrak{g}'^*)^x = 0$
so

$$\mathfrak{g} = \mathbb{R}X + \mathfrak{g}' \quad \text{is } (2,3)\text{-trivial}$$

In this case $\mathcal{P}_{\mathfrak{g}} \cong \wedge^2 \mathfrak{g}'$

Remark:

$$\mathfrak{k} = (0, 0, 12, 13, 23, 14+25+23, \\ 16+25+35+24)$$

has all derivations nilpotent

so $\mathfrak{k} \neq \mathfrak{g}'$ for any $(2,3)$ -trivial $\mathfrak{g}$.

Orbits

If G acts transitively on M then

$$\begin{aligned}\Psi: M &\longrightarrow Z^3(\mathfrak{g}) \\ \langle \Psi, x \wedge y \wedge z \rangle &= c(x, y, z)\end{aligned}\quad (*)$$

has image $G \cdot \Psi_0 = \mathcal{O}'$

$\mathcal{O}'$ defines a strong geometry on $M = G$ via $(*)$

this descends to $M = G/H$ for any H closed in G with
 $\mathfrak{h} \subset \ker \Psi_0 = \{x \in \mathfrak{g} \mid x \cdot \Psi_0 = 0\}$

Suppose $b_2(\mathfrak{g}) = 0$
 then $d\psi: \mathcal{P}_\mathfrak{g}^* \hookrightarrow Z^3(\mathfrak{g})$

For $\beta_0 \in \mathcal{P}_\mathfrak{g}^*$, $\mathcal{O} = G \cdot \beta_0 \xrightarrow{\cong} \mathcal{O}' = G \cdot \Psi_0$
 $\Psi_0 = d\beta_0$
 $\nu(gH) = g \cdot \beta_0$ is the multi-moment map

Proposition: $\mathcal{O} = G \cdot \beta_0 \subset \mathcal{P}_\mathfrak{g}^*$, $b_2(\mathfrak{g}) = 0$, carries a strong geometry in this way only if $\text{stab}_\mathfrak{g} \beta_0 = \ker(d\psi(\beta_0))$

Then $\mathcal{O}$ is 2-plectic and ν is the inclusion $\mathcal{O} \hookrightarrow \mathcal{P}_\mathfrak{g}^*$.

Satisfy $\forall x \in TM|_{\mathcal{O}} \exists p \in \mathcal{P}_\mathfrak{g}^*$ s.t.
 $c(x \wedge p) \neq 0$

MULTI - MOMENT MAPS

III G_2 - geometry

$$\mathbb{R} \subset \mathbb{C} \subset \mathbb{H} \subset \mathbb{O}$$

$$(p, q)(p', q') = (pp' - \bar{q}'q, pq' + \bar{p}'q)$$

$$\overline{(p, q)} = (\bar{p}, -q) \quad \|v\|^2 = v\bar{v}$$

$$\|vw\|^2 = \|v\|^2 \|w\|^2$$

Definition: $G_2 = \text{Aut } \mathbb{O}$

$$V = \mathbb{R}^7 = \text{Im } \mathbb{O}$$

$$\phi_0(x, y, z) = \langle xy, z \rangle, \quad \phi_0 \in \wedge^3 V^*$$

is fully non-degenerate
 $= \langle x \times y, z \rangle$

$$\phi_0 = e_{123} + e_1(e_{45} + e_{67})$$

$$+ e_2(e_{46} + e_{75}) - e_3(e_{47} + e_{56})$$

determines
via $g_0 = \sum e_i^2$ and $\text{vol}_0 = e_{1234567}$

$$(X \lrcorner \phi_0) \wedge (Y \lrcorner \phi_0) \wedge \phi_0 = 6g_0(X, Y) \text{vol}_0$$

$$G_2 = \text{stab}_{GL(7, \mathbb{R})} \phi_0$$

is compact, 14-dimensional, simple
 simply-connected and ϕ_0 stable form

G_2 acts transitively on $S^6 \subset \mathbb{R}^7$

$$e_1 \lrcorner \phi = e_{23} + e_{45} + e_{67}$$

e_1 acts as a complex structure on $e_1^\perp$

$$\phi_0|_{e_1^\perp} = \text{Re}(\text{vol}_e) \Rightarrow e_1^\perp \text{ has an } SU(3) \text{ structure}$$

Definition: $\phi \in \Omega^3(M^7)$ is a G_2 -structure if each $\phi_x \in \Lambda^3 T_x^* M$ is linearly equivalent to ϕ_0 on $\mathbb{R}^7$.

$\Rightarrow M$ oriented, with Riemannian g

(M, ϕ) is torsion-free if $d\phi = 0 = d(*\phi)$

$\Rightarrow \text{Hol}(g) \leq G_2$ & g is Ricci-flat

Suppose T^2 acts preserving ϕ
generators U_1, U_2

Multi-moment $v: M \rightarrow \mathbb{R} \cong \mathcal{D}_g^*$
is just an invariant function with

$$dv = U_1 \lrcorner U_2 \lrcorner \phi \\ = (U_1 \times U_2)^b$$

$$\begin{aligned} d\phi = 0 & \Rightarrow \omega_1 = U_1 \lrcorner \phi, \quad \omega_2 = U_2 \lrcorner \phi \\ d(*\phi) = 0 & \Rightarrow \omega_0 = (U_1 \wedge U_2) \lrcorner * \phi \\ & \text{closed 2-forms} \end{aligned}$$

Proposition:

$$\phi = h^2 \omega_0 \wedge dv + \omega_1 \wedge \Theta_1 + \omega_2 \wedge \Theta_2 \\ + dv \wedge \Theta_1 \wedge \Theta_2$$

where $\Theta_i(U_j) = \delta_{ij}$ $\Theta_i \in \text{span}\{U_1^b, U_2^b\}$

$$\text{and } h^2 = \frac{1}{|U_1 \wedge U_2|^2}$$

For r a regular value of v
 $v^{-1}(r)$ is a hypersurface
 and inherits an $SU(3)$ -structure

$$\sigma = N \lrcorner \phi = h i^* \omega_0 + h^{-1} \theta_1 \wedge \theta_2$$

$$\psi_+ = i^* \phi = i^* \omega_1 \wedge \theta_1 + i^* \omega_2 \wedge \theta_2$$

Proposition: (a) any hypersurface $(Y^6, \sigma, \psi_{\pm})$
 is half-flat

$$\text{i.e. } d(\tfrac{1}{2}\sigma^2) = 0, \quad d\psi_+ = 0$$

(b) a half-flat $(Y^6, \sigma, \psi_{\pm})$ is such a
 hypersurface if and only if one can solve

$$\begin{cases} \frac{d}{dt} \psi_+ = d\sigma \\ \frac{d}{dt} (\tfrac{1}{2}\sigma^2) = -d\psi_- \end{cases}$$

near $Y \times 0 \subset Y \times \mathbb{R}$

the Hitchin
 flow

The modified flow

$$h^{\pm} = \frac{1}{|U_1 \wedge U_2|}$$

$$\begin{cases} \frac{d}{dr} \psi_+ = d(h\sigma) \\ \frac{d}{dr} (\tfrac{1}{2}\sigma^2) = -d(h\psi_-) \end{cases}$$

for $(Y^6, \sigma, \psi_{\pm}, T^2)$ with analytic data
 corresponds to (M^7, ϕ, T^2, v) .

Definition: the reduction of (M, ϕ, T^2, ν)
at level $r \in \mathbb{C}$

$$Z = \nu^{-1}(r)/T^2$$

$$\dim Z = 4$$

$i^*\omega_0, i^*\omega_1, i^*\omega_2$ descend to symplectic

$$\sigma_0, \sigma_1, \sigma_2$$

satisfying

$$\begin{cases} \sigma_0 \wedge \sigma_1 = 0 = \sigma_0 \wedge \sigma_2 \\ h^2 \sigma_0^2 = 2 \text{vol}_Z = \frac{1}{|U_1|^2} \sigma_1^2 = \frac{1}{|U_2|^2} \sigma_2^2 \\ \sigma_1 \wedge \sigma_2 = 2 \langle U_1, U_2 \rangle \text{vol}_Z. \end{cases}$$

Example: $M = \mathbb{R}^7$ $\phi = \phi_0$ $T^2 \subset G_2$ maximal torus

$$\mathbb{R}^7 = \mathbb{R} + \mathbb{C}^3$$

$$(x, z_1, z_2, z_3) \rightarrow (x, e^{i\theta} z_1, e^{i\varphi} z_2, e^{-i(\theta+\varphi)} z_3)$$

$$\nu = -\frac{1}{4} \text{Re}(z_1 z_2 z_3)$$

For $r \neq 0$, $Z = \nu^{-1}(r)/T^2 \equiv \mathbb{R}^4$
diffeo

$$\sigma_0 = \frac{1}{4} (dx \wedge dw + du_2 \wedge du_1)$$

$$\sigma_1 = \frac{1}{2} (dx \wedge du_1 + dw \wedge \eta_2)$$

$$\sigma_2 = \frac{1}{2} (dx \wedge du_2 + \eta_1 \wedge dw)$$

$$\eta_1 = \frac{h^2}{4} (|U_2|^2 du_1 - \langle U_1, U_2 \rangle du_2)$$

$$\eta_2 = \frac{h^2}{4} (|U_1|^2 du_2 - \langle U_1, U_2 \rangle du_1)$$

not hyper Kähler.

Definition: a coherent symplectic triple is $\sigma_0, \sigma_1, \sigma_2$ symplectic forms on Z^4 such that

- (i) $\text{Span} \{ \sigma_0, \sigma_1, \sigma_2 \} \leq \Lambda^2 T^*M$
is a maximal positive subspace
- (ii) $\sigma_0 \wedge \sigma_1 = 0 = \sigma_0 \wedge \sigma_2$

Put $Q = (g_{ij})_{i,j=1}^2$

$$\sigma_i \wedge \sigma_j = g_{ij} \sigma_0^2, \quad h = \sqrt{\det Q}$$

Theorem: Let Y be a principle T^2 -bundle over Z , connection 1-forms θ_1, θ_2 .

Then Y carries an $SU(3)$ -structure

$$\begin{cases} \sigma = h\sigma_0 + h^{-1}\theta_1 \wedge \theta_2 \\ \psi_+ = \sigma_1 \wedge \theta_1 + \sigma_2 \wedge \theta_2 \end{cases}$$

It is half-flat

$$\Leftrightarrow \begin{cases} \sigma_1 \wedge d\theta_1 + \sigma_2 \wedge d\theta_2 = 0 \\ \sigma_0 \wedge d\theta_1 = 0 = \sigma_0 \wedge d\theta_2 \end{cases}$$

If the data is analytic then it defines a torsion-free G_2 -structure with T^2 -symmetry on a neighbourhood of $Y \times 0$ in $Y \times \mathbb{R}$

Every simply-connected torsion-free G_2 -structure with free T^2 -symmetry arises in this way.

The flow may be described by the following equations Z :

$$\sigma'_0 = 0$$

$$\sigma'_1 = -d\theta_2, \quad \sigma'_2 = d\theta_1$$

$$g'_{11} \sigma_0^2 = -2\sigma_1 \wedge d\theta_2$$

$$g'_{22} \sigma_0^2 = 2\sigma_2 \wedge d\theta_1$$

$$g'_{12} \sigma_0^2 = \sigma_1 \wedge d\theta_1 + \sigma_2 \wedge d\theta_2$$

$$\sigma_0 \wedge \theta'_1 = dg_{12} \wedge \sigma_2 - dg_{22} \wedge \sigma_1$$

$$\sigma_0 \wedge \theta'_2 = dg_{11} \wedge \sigma_2 - dg_{12} \wedge \sigma_1$$

Example: Z hyperKähler

$$\sigma_0^2 = \sigma_1^2 = \sigma_2^2, \quad \sigma_i \wedge \sigma_j = 0 \quad (i \neq j)$$

$$[\sigma_1], [\sigma_2] \in H^2(\mathbb{Z}, \mathbb{Z})$$

$$d\theta_1 = a\sigma_1 + b\sigma_2$$

$$a, b \in \mathbb{Z}$$

$$d\theta_2 = a\sigma_1 - a\sigma_2$$

gives half-complete G_2 -metrics
for $\det \begin{vmatrix} a & b \\ a & -a \end{vmatrix} \gg 0$.

$d\theta_i \notin \text{span} \{\sigma_0, \sigma_1, \sigma_2\}$ is more
interesting

Singular behaviour?

Nearly Kähler, Spin(7), ...